

SEARCHING FOR SPECTROSCOPIC BINARIES WITHIN TRANSITION DISK OBJECTS¹

Saul A. Kohn

Lowell Observatory, 1400 W. Mars Hill Road, Flagstaff, AZ 86001, USA

*Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, PA,
19104, USA*

`saulkohn@sas.upenn.edu`

Evgenya L. Shkolnik

School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA

Lowell Observatory, 1400 W. Mars Hill Road, Flagstaff, AZ 86001, USA

`shkolnik@asu.edu`

Alycia J. Weinberger

*Department of Terrestrial Magnetism, Carnegie Institution of Washington, 5241 Broad
Branch Road, NW, Washington, DC 20015, USA*

`weinberger@dtm.ciw.edu`

Joleen K. Carlberg

NASA/GSFC Code 667, Greenbelt, MD 20771, USA

`joleen.k.carlberg@nasa.gov`

and

Joe Llama

Lowell Observatory, 1400 W. Mars Hill Road, Flagstaff, AZ 86001, USA

SUPA, School of Physics & Astronomy, North Haugh. St Andrews. Fife. KY16 9SS, UK

`joe.llama@st-andrews.ac.uk`

ABSTRACT

Transition disks (TDs) are intermediate stage circumstellar disks characterized by an inner gap within the disk structure. To test whether these gaps may have been formed by closely orbiting, previously undetected stellar companions, we collected high-resolution optical spectra of 31 TD objects to search for spectroscopic binaries (SBs). Twenty-four of these objects are in Ophiuchus and seven are within the Coronet, Corona Australis, and Chameleon I star-forming regions. We measured radial velocities for multiple epochs, obtaining a median precision of 400 ms^{-1} . We identified double-lined SB SSTc2d J163154.7-250324 in Ophiuchus, which we determined to be composed of a K7(± 0.5) and a K9(± 0.5) star, with orbital limits of $a < 0.6 \text{ AU}$ and $P < 150 \text{ days}$. This results in an SB fraction of $0.04^{+0.12}_{-0.03}$ in Ophiuchus, which is consistent with other spectroscopic surveys of non-TD objects in the region. This similarity suggests that TDs are not preferentially sculpted by the presence of close binaries and that planet formation around close binaries may take place over similar timescales to that around single stars.

Subject headings: Binaries: spectroscopic – Circumstellar matter – Optical: stars
– Stars: pre-main sequence

1. Introduction

Transition disks (TDs) are an intermediate stage of circumstellar disk evolution among pre-main sequence (PMS) stars. They are characterized by a lack of near-infrared excesses caused by inner optically thin holes or ‘gaps’ opening within the disk structure, which typically have radii between one and tens of AU, and large far-infrared excesses from the disk material at larger separations (e.g., Hughes et al. 2007, 2009; Andrews et al. 2011; Zhu et al. 2011).

Planet formation and dust aggregation are thought to cause the gaps observed in TDs, with inner holes greater than 15 AU possibly due to the formation of multiple planets (Dodson-Robinson & Salyk 2011; Zhu et al. 2011). Other gap-formation mechanisms may involve ultraviolet (UV) and X-ray photons irradiating the disk and accretion from the disk onto the host star (e.g. Gorti et al. 2009; Espaillat et al. 2012). Recent observations by Follette et al. (2013) support the results of simulations from, e.g. Rice et al. (2006) and Zhu

¹This paper is based on data gathered with the 6.5 m Clay Telescope located at Las Campanas Observatory, Chile.

& Stone (2014) that dust dynamics may play a role in gap formation. The distribution of large and small grains in the inner regions of disks can be modified by gas pressure gradients, perhaps also generated by planets, that may trap or filter large grains as well as allow grains to grow efficiently (e.g. Zhu et al. 2011; Pinilla et al. 2012).

A closely orbiting stellar companion may also produce an inner gap in a circumstellar disk (Artymowicz & Lubow 1994; D’Alessio et al. 2005), as illustrated by the CoKu Tauri/4 system. This system was observed to have a 10 AU wide gap (Forrest et al. 2004). The initial interpretation for such a wide gap was that a $10 M_{\text{Jup}}$ planet orbiting the central star cleared out the material (Quillen et al. 2004). However, it was afterward shown that CoKu Tauri/4 was in fact a close visual binary (VB) system with a separation of ~ 8 AU (Ireland & Kraus 2008), and that the secondary star likely carved out the gap. Similarly, Biller et al. (2012) also detected a VB within the TD object HD 142527. With a projected separation of $(12.8 \pm 1.5 \text{ AU})$, the companion is likely to have played a role in shaping the system’s complex disk structure, and it may have been responsible for the disk’s cleared inner ring.

A close binary system need not preclude circumbinary planet formation. Using photometric data from *Kepler*, Doyle et al. (2011) announced the first discovery of a circumbinary planet. Several other circumbinary planets have since been discovered (e.g. Orosz et al. 2012; Qian et al. 2012; Qian et al. 2012b; Schwamb et al. 2013). These examples illustrate the need to further study multiple-star systems that host circumbinary material, and to test the effects of companion stars on disk dissipation and planet formation. Young star-forming regions are good test beds for such a study.

The nearby (140 pc) Ophiuchus star-forming region (Oph SFR) contains a population of young stars of an average age ~ 2.2 Myr (Willing et al. 2005; M. Lombardi & Alves 2008; Erickson et al. 2011). More than 50% of the young stellar objects (YSOs) in Oph show evidence of optically thick circumstellar disks at mid-IR wavelengths (Bontemps et al. 2001), of which 9% are identified as TDs (Evans et al. 2003; Cieza et al. 2010). This makes it well-suited for studies of disk dissipation.

Previous surveys of multiplicity in Oph and other SFRs have concentrated on visual binarity across all of the known active star-forming regions (e.g. Table 6 of Lafreniere et al. 2008). With the exception of interferometric methods that are sensitive to separations of order 0.1–1 AU (Pott et al. 2010), VB searches are typically limited to separations > 4 AU (Ghez et al. 1993; Kraus et al. 2012). Cieza et al. (2010) (hereafter C10) used IR and optical data to identify the TDs in Oph. C10 also sub-classify the TDs in their sample into five categories: grain growth-dominated disks (which are accreting and have a negative slope of IR excess: 44% of their sample), giant planet-forming disks (accreting with a positive slope of IR excess: 13%), photoevaporating disks (non-accreting with high disk luminosity: 17%),

disks already in the debris disk stage (non-accreting with low disk luminosity: 13%) and circumbinary disks (13%). C10 did not observe an increase in detected companions within the separations to which their survey was most sensitive to (8–20 AU) compared with other separations, leading them to suggest that stellar companions at these separations are not responsible for a large fraction of the TD population.

We present a complementary radial velocity (RV) survey to C10, to search for stellar companions closer in to the primary star. Our aim is to test whether the spectroscopic binary (SB) fraction is similar or enhanced in the TD population relative to stars without such disks. The structure of this paper is as follows. The selection of the TD systems, observations and data reduction are described in Section 2. RV, rotational velocity, and spectral type (SpT) measurements are presented in Section 3. We report the discovery of a double-lined spectroscopic binary in Section 4, and evaluate possible biases and the completeness of our survey in Section 5. We derive a multiplicity fraction for Oph TDs, compare with previous multiplicity surveys of Oph, and discuss what multiplicity among TD objects might mean for theories of disk dispersal in Section 6. We summarize and conclude our study in Section 7. While not the main objective of this paper, we provide values of lithium I (Li I) and H α equivalent widths (EWs), and accretion rates for our sample, and a discussion on the unusual nature of Oph star SR21A in the Appendix.

2. Sample and Observations

To survey the innermost regions of TD objects for stellar companions, we collected multi-epoch, high-resolution optical spectroscopy of 24 stars in Oph identified as having TDs (C10; Andrews & Williams 2007a; Geers et al. 2007; Furlan et al. 2009), two in Corona Australis (CrA; <3 Myr; Neuhauser & Forbrich 2008; Sicilia-Aguilar et al. 2008; Hughes et al. 2010), two in Coronet (Cor, embedded within CrA; <1 Myr; Knacke et al. 1973; Sicilia-Aguilar et al. 2008; Hughes et al. 2010), and three in Chameleon I (Cha; 1–3 Myr; Furlan et al. 2009; Kim et al. 2009).

C10 identified Oph TDs from their colors as measured by the *Spitzer Cores 2 Disks* legacy program (Evans et al. 2003), and 20 of our 24 targets come from that work. A *Spitzer* color (where $[x]$ indicates the magnitude in the $x \mu\text{m}$ band) of $[3.6] - [4.5] < 0.25$ indicates the flux deficit characteristic of the near-infrared/inner opacity holes of a TD, while a color of $[3.6] - [24] > 1.5$ ensures that all of the targets have significant mid-IR excesses at $24 \mu\text{m}$. We included an additional four Oph targets and seven from other regions, all of which were identified by other authors using the same criteria as C10 (Andrews & Williams 2007a; Geers et al. 2007; Furlan et al. 2009).

The Oph targets in this study are shown in Figure 1 in relation to other surveys of YSOs in the region. Table 1 lists each target in our sample, its stellar association, coordinates, SpT, and V , R and J magnitudes. Targets that have known wide stellar companions are also noted. The SpT distribution of our sample is shown in Figure 2 and is dominated by late-K to mid-M stars with a median SpT of M2.

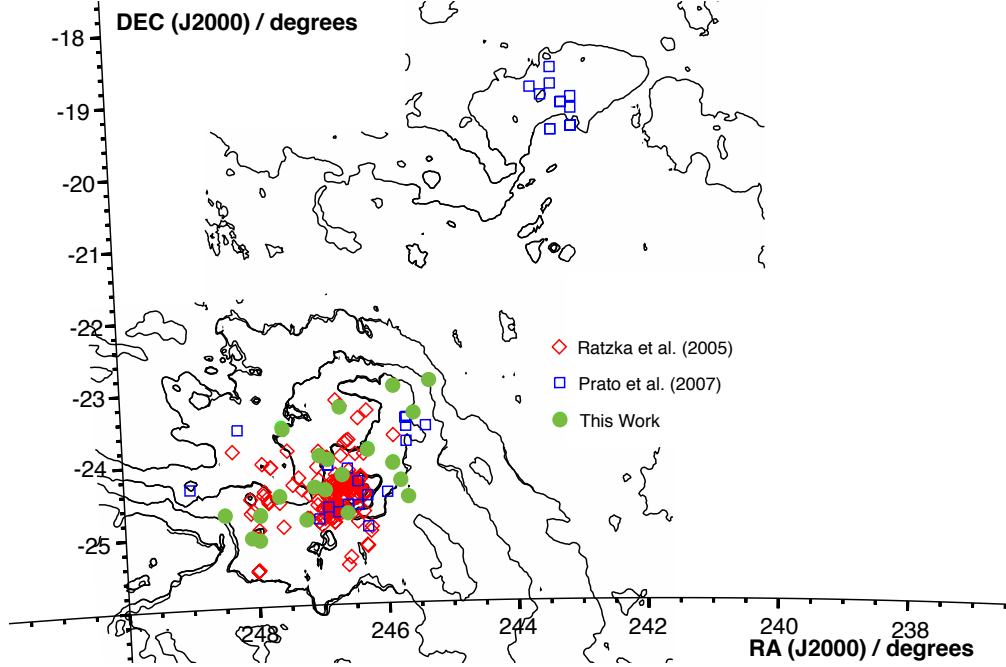


Fig. 1.— R.A. and Dec. of our targets within Ophiuchus, overlaid with other multiplicity surveys. Contours are from the *Infrared Astronomical Satellite* (IRAS) *Sky Survey Atlas* at $100\mu\text{m}$ (Neugebauer et al. 1984; Beichman et al. 1988).

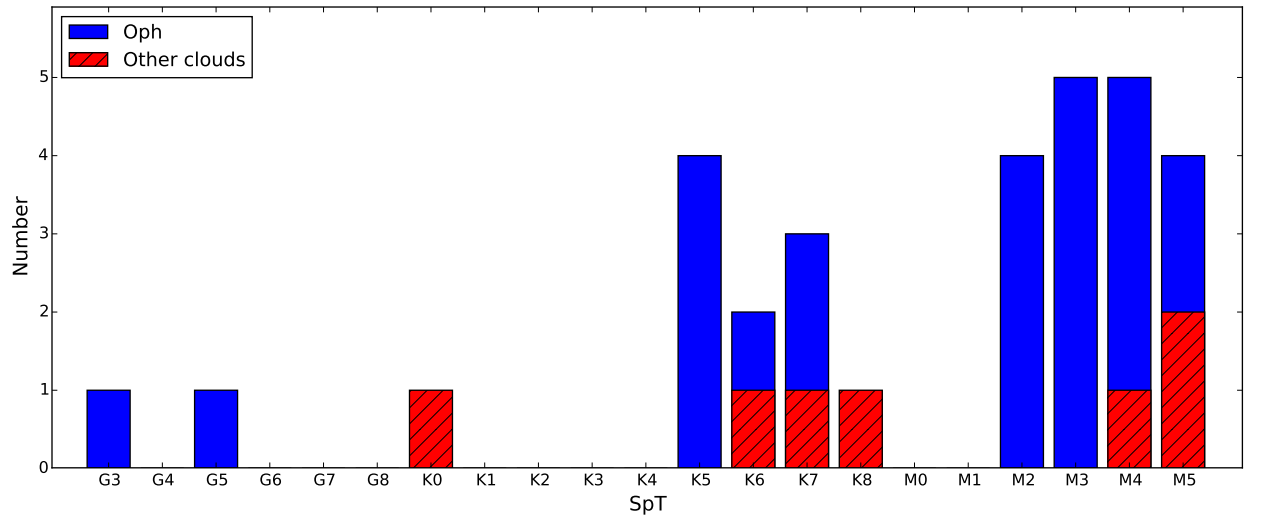


Fig. 2.— The spectral type distribution of the TD sample. All of the targets have ages ≤ 10 Myr.

Table 1. Transition Disk Targets

Name	Other Identifier	Association	α (J2000) h m s	δ (J2000) ° ' "	SpT	V mag	R mag	J (2MASS) mag	Known Visual Multiplicity ^a (Separation in ")	Source ^b
SZ Cha		Cha	10 58 16.8	-77 17 17.1	K0	12.68		9.25	Binary (~5)	(1), (2), (3), (4)
T25	Sz 18	Cha	11 07 19.1	-76 03 04.8	M2.5	15.35	13.7	10.96		(1), (5)
T35	Sz 27	Cha	11 08 39.0	-77 16 04.2	K8		15.82	11.17		(1), (5)
RX J1852.3-3700		Cha	18 52 17.3	-37 00 12.0	K7	12.19		9.77		(7),(8)
CrA-4111		CrA	19 01 20.8	-37 03 03.0	M4.5			13.23		(9)
G-49	CrA 468	Cor	19 01 49.4	-37 00 28.0	M4		16.6	12.5		(9), (10)
G-102	CrA 133	Cor	19 01 25.6	-37 04 53.0	M5		15.3	12.36		(9), (5)
SSTc2d J162118.5-225458		Oph	16 21 18.5	-22 54 58.0	M2			11.45		(11)
SSTc2d J162218.5-232148		Oph	16 22 18.5	-23 21 48.0	K5	12.67		9.52		(11), (12)
SSTc2d J162245.4-243124		Oph	16 22 45.4	-24 31 24.0	M3		14.2	10.38		(11), (5)
SSTc2d J162309.2-241705		Oph	16 23 09.3	-24 17 03.0	M?	12.75	14.2	10.32		(1), (5)
SSTc2d J162332.8-225847		Oph	16 23 32.8	-22 58 47.0	M5		15.7	11.49		(11), (5)
SSTc2d J162336.1-240221		Oph	16 23 36.1	-24 2 21.0	M5	14.68		11.53		(11), (13)
SSTc2d J162506.9-235050		Oph	16 25 06.9	-23 50 50.0	M3			11.05		(11)
SSTc2d J162623.7-244314	DoAr 25	Oph	16 26 23.7	-24 43 14.0	K5		12.65	9.39		(11), (14)
SSTc2d J162646.4-241160		Oph	16 26 46.4	-24 11 60.0	G5		13.9	9.68	Binary (0.58)	(11), (14), (15)
DoAr28	Haro 1-8	Oph	16 26 47.4	-23 14 52.2	K?	13.84	12.1	9.89		(16), (17), (5)
SR21A		Oph	16 27 10.3	-24 19 12.7	G3	14.1		8.75	Binary (896)	(18), (19), (20)
SSTc2d J162738.3-235732	DoAr 32	Oph	16 27 38.3	-23 57 32.0	K5		13.66	9.91	Triple(?) ^c	(11), (14), (15), (21)
SSTc2d J162739.0-235818	DoAr 33	Oph	16 27 39.0	-23 58 18.0	K6		13.24	9.9	Triple(?) ^c	(11), (14), (15), (21)
SSTc2d J162740.3-242204	DoAr 34	Oph	16 27 40.3	-24 22 04.0	K5	11.5	11.1	8.44		(11), (22)
SSTc2d J162802.6-235504		Oph	16 28 02.6	-23 55 04.0	M3			11.76		(11)
SSTc2d J162821.5-242155		Oph	16 28 21.5	-24 21 55.0	M3		16.97	12.08		(11), (14)
SSTc2d J162854.1-244744		Oph	16 28 54.1	-24 47 44.0	M2		15.24	10.68		(11), (14)
SSTc2d J163020.0-233108		Oph	16 30 20.0	-23 31 08.0	M4			11.32		(11)
SSTc2d J163033.9-242806		Oph	16 30 33.9	-24 28 06.0	M4			11.63		(11)
SSTc2d J163145.4-244307		Oph	16 31 45.4	-24 43 07.0	M4			11.8		(11)
SSTc2d J163154.4-250349		Oph	16 31 54.4	-25 03 49.0	M4			11.78		(11)
SSTc2d J163154.7-250324	Ha74	Oph	16 31 54.7	-25 03 24.0	K7	12.64	13.3	10.14		(11), (5)
SSTc2d J163205.5-250236		Oph	16 32 05.5	-25 02 36.0	M2		15.5	11.65		(11), (5)
SSTc2d J163355.6-244205		Oph	16 33 55.6	-24 42 05.0	K7		14.1	10.46		(11), (5)

^aBinary companions all lie outside of the TD.

^bReferences: (1) Furlan et al. (2009), (2) Kim et al. (2009), (3) Ducati (2002), (4) Vogt et al. (2012), (5) Cutri et al. (2003), (6) Lopez Martí et al. (2013), (7) Hughes et al. (2010), (8) Kiraga (2012), (9) Sicilia-Aguilar et al. (2006), (10) Monet et al. (2003), (11) Cieza et al. (2010), (12) Vrba et al. (1993), (13) Allers et al. (2006), (14) Wilking et al. (2005), (15) Ratzka et al. (2005), (16) Andrews & Williams (2007a), (17) Samus' et al. (2003), (18) Geers et al. (2007), (19) Richichi & Percheron (2002), (20) Prato et al. (2003), (21) Barsony & Marsh (2005). The SpTs are from the first reference listed, followed by the references for V and R magnitudes, and then visual multiplicity, if applicable. All J magnitudes are from the 2MASS catalog (Skrutskie et al. 2006).

^cSSTc2d J162738.3-235732 and SSTc2d J162739.0-235818 are both possible members of the same triple system about primary ROXs 30A (16h 27m 37.0s -23°59'32''; Ratzka et al. 2005).

2.1. Observations

We acquired high-resolution optical spectra using the Magellan Inamori Kyocera Echelle (MIKE; Bernstein et al. 2003) spectrograph at the Magellan (Clay) 6.5 m telescope. Twenty-one of our targets were observed twice, two or three days apart during 18–23 June, 2010 (UT100618–UT100623)², and then nearly a year later on 14–15 June, 2011 (UT110614–UT110615). Two targets were observed a third time on 9–10 May, 2012 (UT120509–UT120510), and another on 17 January, 2012 (UT120117). The complete log of our observations is recorded in Table 2. We used the red arm of the spectrograph and the 0.5'' slit to obtain a spectrum from 4900 to 9150 Å with a spectral resolution of $\sim 47,000$. Total exposure times ranged from 2 to 30 minutes per target depending on stellar brightness to achieve a typical signal-to-noise ratio (S/N) of ≈ 50 per resolution element at 8000 Å. For the faintest stars, the S/N was roughly 10–20.

Data were reduced using the facility pipeline (Kelson 2003). Each stellar exposure was bias-subtracted and flat-fielded for pixel-to-pixel sensitivity variations. After optimal extraction, spectra were wavelength calibrated with a ThAr arc taken within an hour of the stellar exposure. To correct for instrumental drift, the telluric molecular oxygen A band (7619 – 7657 Å) was used to align the MIKE spectra to an average precision of 70 m s^{−1}. We then corrected for the heliocentric velocity.

3. Results

3.1. Radial Velocities

Each target spectrum was cross-correlated against an RV standard star best matched to its SpT and observed on the same night. We used the IRAF³ *fxcor* routine (Wyatt 1985; Fitzpatrick 1993) to obtain RVs and search for single- and double-lined spectroscopic binaries (SB1s and SB2s, respectively). The standards and their published RVs are listed in Table 3. We used orders 39, 40, 43, 44, 46, 48, 49, 51, and 53, which ranged from 6400 to 8900 Å, for the cross-correlation to avoid telluric absorption lines and very low S/N regions. The RVs reported in Table 2 are the averages of those measured for each order weighted by the

²Throughout this paper we use a date format of UTYYMMDD.

³IRAF is distributed by the National Optical Astronomy Observatories, which are operated by the Association of Universities for Research in Astronomy, Inc., under a cooperative agreement with the National Science Foundation.

average S/N in each order. Table 2 also lists the standard star used, and Li I and H α EWs (see Appendix 8.1 for details of Li I and H α analysis).

The uncertainties in our RV measurements include the standard error on the mean RV across the cross-correlated apertures (typically $\sim 400 \text{ ms}^{-1}$), an uncertainty of 300 ms^{-1} to account for RV distortions due to star spots on young stars (Hatzes 2002; Mohanty et al. 2002; Berdyugina 2005; Desort et al. 2007), and an additional 400 ms^{-1} to account for the zero-point uncertainty in the RVs of our standard stars (see Table 3), all added in quadrature. In all cases, the systematic uncertainties from possible spot activity and the zero-point dominate the final uncertainties. Note that $\Delta(\text{RV})$, the difference in RVs between epochs, need not include the uncertainty in the RV of the standards as the same standards were used.

The largest systematic uncertainties in these results will come from mismatched SpTs between the TD objects and the RV standard stars. While we endeavoured to match SpTs as closely as possible, the lack of G or K stars among our RV standards implies that for G and K-type TDs, systematics may remain unaccounted for. However, this should only affect SZ Cha (K0), SSTc2d SSTc2d J162646.4-241160 (G5) and SR21A (G3–K?; a discussion of its SpT can be found in Appendix 8.2), since late K-type stars can be well matched by early M-type standards. Indeed, we see that the uncertainties on the RVs of SZ Cha are larger ($\sim 2\times$) than the average uncertainties.

The average RV of our Oph targets is -6.6 kms^{-1} with a velocity dispersion of 1.3 kms^{-1} , which agrees well with other measurements for the RV of the cloud: $-5.64 \pm 2.57 \text{ kms}^{-1}$ from Kurosawa et al. (2006) and $-6.27 \pm 1.48 \text{ kms}^{-1}$ from Prato (2007). There are no overlapping targets between our sample and those studies.

Our average RV for Cha targets is $15 \pm 1 \text{ kms}^{-1}$, which is in good agreement with James et al. (2006), who report an average of $12.8 \pm 3.6 \text{ kms}^{-1}$. However, our measurements of the average RV of CrA+Cor, $-3 \pm 1 \text{ kms}^{-1}$, differs by $\sim 1.5\sigma$ from James et al.’s findings of $-1.1 \pm 0.5 \text{ kms}^{-1}$. It is, however, in good agreement with Neuhäuser et al. (2000), who find an average RV of $-2.6 \pm 1.4 \text{ kms}^{-1}$ for ROSAT-selected T Tauri stars in CrA.

The close agreement in mean RV values within our sample implies that any single RV measurement that deviates by more than 3σ from the region average would be a highly probable RV-variable SB1 (Kurosawa et al. 2006; Prato 2007). We find no such targets in our sample.

Table 2. Observation log and measured properties of TD objects

Target	Date observed dd/mm/yyyy	Standard Star Used	RV kms ⁻¹	Li I EW ^a Å	H α EW ^a Å	H α 10% ^b kms ⁻¹
SZ Cha ^c	18/06/2010	GJ 908	12.1 \pm 0.5	0.36 \pm 0.04	-3.68 \pm 0.03	270 \pm 25
	21/06/2010	GJ 908	14.0 \pm 1.5	0.38 \pm 0.02	-26.6 \pm 0.5	479 \pm 47
	14/06/2011	GJ 514	17.0 \pm 1.3	0.42 \pm 0.05	-15.7 \pm 0.9	478 \pm 43
	mean		14.0 \pm 2.0	0.39 \pm 0.07	-14 \pm 2	409 \pm 69
T25 ^c	18/06/2010	GJ 908	15.3 \pm 0.3	0.55 \pm 0.02	-10.7 \pm 0.8	327 \pm 29
	21/06/2010	GJ 908	15.6 \pm 0.3	0.58 \pm 0.02	-11 \pm 2	347 \pm 34
	14/06/2011	GJ 514	16.0 \pm 0.3	0.60 \pm 0.04	-21 \pm 1	444 \pm 49
	mean		15.7 \pm 0.6	0.58 \pm 0.05	-15 \pm 1	373 \pm 66
T35	21/06/2010	GJ 908	16.2 \pm 0.6	0.49 \pm 0.03	-118 \pm 15	482 \pm 45
	14/06/2011	GJ 699	16.8 \pm 0.6	0.41 \pm 0.08	-128 \pm 16	520 \pm 55
	mean		16.5 \pm 0.8	0.45 \pm 0.09	-123 \pm 22	501 \pm 71
RX J1852.3-3700 ^d	19/06/2010	GJ 908	-1.5 \pm 0.3	0.5 \pm 0.1	-45 \pm 6	320 \pm 30
	22/06/2010	GJ 699	-2.6 \pm 0.5	0.5 \pm 0.1	-33 \pm 2	325 \pm 33
	14/06/2011	GJ 514	-3.0 \pm 0.3	0.48 \pm 0.01	-27 \pm 3	326 \pm 35
	mean		-2.4 \pm 0.6	0.5 \pm 0.2	-35 \pm 7	324 \pm 57
CrA-4111	23/06/2010	GJ 699	-5.2 \pm 0.3	< 0.16	-16 \pm 4	135 \pm 14
G-49	15/06/2011	GJ 699	-2.7 \pm 0.3	0.51 \pm 0.04	-3 \pm 2	92 \pm 9
	09/05/2012	GJ 699	-2.6 \pm 0.3	0.62 \pm 0.01	-4.32 \pm 0.05	106 \pm 10
	10/05/2012	GJ 699	-2.8 \pm 0.3	0.62 \pm 0.01	-4.71 \pm 0.02	101 \pm 9
	mean		-2.7 \pm 0.6	0.55 \pm 0.09	-4 \pm 2	100 \pm 16
G-102	15/06/2011	GJ 699	-2.8 \pm 0.5	0.61 \pm 0.03	-14 \pm 4	189 \pm 18
	09/05/2012	GJ 699	-2.9 \pm 0.3	0.66 \pm 0.01	-20 \pm 2	153 \pm 14
	10/05/2012	GJ 699	-1.0 \pm 0.3	0.67 \pm 0.01	-17 \pm 1	152 \pm 7
	mean		-2.2 \pm 0.7	0.65 \pm 0.03	-17 \pm 4	165 \pm 24
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SSTc2d J162118.5-225458 ^c	21/06/2010	GJ 908	-6.3 \pm 0.4	0.48 \pm 0.01	-14.3 \pm 0.4	336 \pm 36
	22/06/2010	GJ 908	-6.6 \pm 0.4	0.5 \pm 0.01	-8.9 \pm 0.7	285 \pm 30
	mean		-6.4 \pm 0.5	0.49 \pm 0.01	-11.6 \pm 0.8	310 \pm 47
SSTc2d J162218.5-232148 ^c	19/06/2010	GJ 908	-7.9 \pm 0.4	0.4 \pm 0.2	-8.2 \pm 0.2	362 \pm 40
	21/06/2010	GJ 908	-9.7 \pm 0.4	0.4 \pm 0.1	-15 \pm 1	551 \pm 53
	14/06/2011	GJ 514	-7.5 \pm 0.4	0.42 \pm 0.03	-12.5 \pm 0.7	378 \pm 37
	mean		-8.4 \pm 0.7	0.4 \pm 0.2	-12 \pm 2	430 \pm 76
SSTc2d J162245.4-243124	20/06/2010	GJ 699	-5.3 \pm 0.3	0.5 \pm 0.4	-4.6 \pm 0.4	120 \pm 13
	22/06/2010	GJ 699	-5.4 \pm 0.3	0.51 \pm 0.01	-5.1 \pm 0.8	124 \pm 12
	14/06/2011	GJ 699	-4.0 \pm 0.3	0.55 \pm 0.05	-4.2 \pm 0.8	115 \pm 4
	mean		-4.8 \pm 0.6	0.5 \pm 0.4	-5 \pm 1	120 \pm 18
SSTc2d J162309.2-241705	20/06/2010	GJ 908	-3.8 \pm 0.4	0.28 \pm 0.06	-17.3 \pm 0.5	419 \pm 46
	22/06/2010	GJ 908	-2.9 \pm 0.4	0.29 \pm 0.01	-16.3 \pm 0.8	365 \pm 30
	14/06/2011	GJ 699	-3.6 \pm 0.4	0.29 \pm 0.01	-14 \pm 1	363 \pm 35
	mean		-3.4 \pm 0.7	0.29 \pm 0.06	-16 \pm 2	382 \pm 65
SSTc2d J162332.8-225847	20/06/2010	GJ 699	-6.9 \pm 0.3	0.54 \pm 0.01	-13 \pm 3	153 \pm 15
	22/06/2010	GJ 699	-7.0 \pm 0.3	0.44 \pm 0.03	-11 \pm 3	122 \pm 13
	mean		-6.9 \pm 0.4	0.49 \pm 0.04	-12 \pm 4	137 \pm 20
SSTc2d J162336.1-240221	20/06/2010	GJ 699	-6.7 \pm 0.3	< 0.28	-12 \pm 3	93 \pm 8

Table 2—Continued

Target	Date observed dd/mm/yyyy	Standard Star Used	RV kms ⁻¹	Li I EW ^a Å	H α EW ^a Å	H α 10% ^b kms ⁻¹
SSTc2d J162506.9-235050 ^c	22/06/2010	GJ 699	-7.3 \pm 0.3	0.36 \pm 0.01	-8 \pm 1	113 \pm 11
	mean		-7.0 \pm 0.4	0.4 \pm 0.3	-10 \pm 3	103 \pm 14
	20/06/2010	GJ 699	-6.5 \pm 0.3	0.51 \pm 0.01	-9 \pm 1	259 \pm 26
	22/06/2010	GJ 699	-6.9 \pm 0.3	0.53 \pm 0.01	-12 \pm 3	341 \pm 34
	14/06/2011	GJ 699	-7.0 \pm 0.3	0.57 \pm 0.01	-14 \pm 4	383 \pm 48
SSTc2d J162623.7-244314	mean		-6.8 \pm 0.6	0.53 \pm 0.01	-12 \pm 5	328 \pm 64
	21/06/2012	GJ 908	-7.2 \pm 0.4	0.5 \pm 0.2	-4.5 \pm 0.4	239 \pm 25
	14/06/2011	GJ 908	-6.7 \pm 0.4	0.56 \pm 0.02	-11 \pm 1	467 \pm 47
	mean		-6.9 \pm 0.6	0.5 \pm 0.2	-8 \pm 1	353 \pm 53
	20/06/2010	GJ 908	-6.3 \pm 0.4	0.5 \pm 0.1	-11.7 \pm 0.4	346 \pm 31
SSTc2d J162646.4-241160	22/06/2010	GJ 908	-6.6 \pm 0.3	0.4 \pm 0.2	-14.7 \pm 0.8	368 \pm 36
	14/06/2011	GJ 514	-6.2 \pm 0.4	0.4 \pm 0.1	-17 \pm 2	356 \pm 34
	mean		-6.4 \pm 0.7	0.4 \pm 0.3	-14.38 \pm 2.2	357 \pm 58
	20/06/2010	GJ 908	-8.6 \pm 0.3	0.45 \pm 0.01	-25 \pm 2	451 \pm 40
	22/06/2010	GJ 908	-8.0 \pm 0.4	0.4 \pm 0.2	-29 \pm 2	441 \pm 47
DoAr28	14/06/2011	GJ 514	-7.9 \pm 0.3	0.4 \pm 0.2	-28 \pm 1	436 \pm 40
	mean		-8.2 \pm 0.6	0.4 \pm 0.3	-27 \pm 3	443 \pm 74
	14/06/2011	GJ 908	-3.6 \pm 0.5	0.13 \pm 0.02	1.1 \pm 0.15 ^f	...
	19/06/2010	GJ 908	-7.4 \pm 0.3	0.47 \pm 0.01	-16 \pm 2	291 \pm 33
	22/06/2010	GJ 908	-6.5 \pm 0.3	0.49 \pm 0.02	-14 \pm 3	275 \pm 29
SSTc2d J162738.3-235732	14/06/2011	GJ 514	-6.4 \pm 0.3	0.5 \pm 0.05	-13.4 \pm 0.9	292 \pm 29
	mean		-6.8 \pm 0.6	0.49 \pm 0.05	-15 \pm 4	286 \pm 53
	19/06/2010	GJ 699	-7.7 \pm 0.5	0.49 \pm 0.01	-22 \pm 1	373 \pm 37
	21/06/2010	GJ 699	-7.3 \pm 0.7	0.52 \pm 0.01	-13 \pm 4	282 \pm 28
	14/06/2011	GJ 514	-7.0 \pm 0.4	0.54 \pm 0.01	-9.9 \pm 0.2	261 \pm 27
SSTc2d J162739.0-235818 ^c	mean		-7.3 \pm 0.9	0.51 \pm 0.02	-15 \pm 4	306 \pm 54
	22/06/2010	GJ 908	-6.0 \pm 0.4	0.5 \pm 0.1	-10.3 \pm 0.2	289 \pm 29
	14/06/2011	GJ 514	-7.2 \pm 0.3	0.51 \pm 0.01	-10.77 \pm 0.09	362 \pm 32
	mean		-6.6 \pm 0.5	0.5 \pm 0.1	-10.5 \pm 0.2	325 \pm 43
	20/06/2010	GJ 908	-7.3 \pm 0.4	0.51 \pm 0.04	< 1	...
SSTc2d J162802.6-235504	22/06/2010	GJ 908	-6.6 \pm 0.3	0.57 \pm 0.08	-3.7 \pm 0.2	77 \pm 6
	mean		-6.9 \pm 0.5	0.54 \pm 0.09	-3.5 \pm 0.9	77 \pm 6
	21/06/2010	GJ 699	-7.1 \pm 0.3	0.4 \pm 0.2	-2.9 \pm 0.6	82 \pm 7
	22/06/2010	GJ 699	-7.5 \pm 0.3	0.4 \pm 0.2	-3 \pm 1	95 \pm 5
	mean		-7.3 \pm 0.5	0.4 \pm 0.3	-3 \pm 1	88 \pm 9
SSTc2d J162854.1-244744	20/06/2010	GJ 908	-5.0 \pm 0.3	0.58 \pm 0.02	-21.7 \pm 0.4	401 \pm 37
	22/06/2010	GJ 908	-5.1 \pm 0.3	0.59 \pm 0.02	-24 \pm 6	314 \pm 29
	14/06/2011	GJ 514	-4.4 \pm 0.3	0.6 \pm 0.3	-65 \pm 13	423 \pm 52
	mean		-4.8 \pm 0.6	0.6 \pm 0.3	-37 \pm 14	379 \pm 70
	21/06/2010	GJ 908	-7.2 \pm 0.4	0.57 \pm 0.75	-2.1 \pm 0.4	87 \pm 7
SSTc2d J163020.0-233108	22/06/2010	GJ 699	-6.7 \pm 0.3	0.56 \pm 0.57	-1.6 \pm 0.4	83 \pm 7
	mean		-7.0 \pm 0.5	0.57 \pm 0.94	-1.8 \pm 0.6	85 \pm 10
	20/06/2010	GJ 699	-6.5 \pm 0.5	1.3 \pm 0.2	-19 \pm 3	332 \pm 33
	22/06/2010	GJ 699	-6.1 \pm 0.4	0.51 \pm 0.08	-17 \pm 1	247 \pm 24
	mean		-6.3 \pm 0.6	0.9 \pm 0.2	-18 \pm 4	290 \pm 41
SSTc2d J163033.9-242806						

Table 2—Continued

Target	Date observed dd/mm/yyyy	Standard Star Used	RV kms ⁻¹	Li I EW ^a Å	H α EW ^a Å	H α 10% ^b kms ⁻¹
SSTc2d J163145.4-244307	21/06/2010	GJ 699	-5.0 \pm 0.4	0.5 \pm 0.4	-12.2 \pm 0.2	455 \pm 49
	22/06/2010	GJ 699	-2.4 \pm 0.4	0.3 \pm 0.2	-37 \pm 2	446 \pm 39
	14/06/2011	GJ 699	-4.8 \pm 0.3	0.5 \pm 0.05	-60 \pm 9	364 \pm 37
	mean		-4.1 \pm 0.7	0.4 \pm 0.4	-36 \pm 9	422 \pm 73
SSTc2d J163154.4-250349	21/06/2010	GJ 699	-4.6 \pm 0.4	0.3 \pm 0.1	-106 \pm 29	374 \pm 43
	14/06/2011	GJ 514	-3.7 \pm 0.3	0.34 \pm 0.07	-139 \pm 33	402 \pm 35
	mean		-4.1 \pm 0.5	0.3 \pm 0.1	-123 \pm 44	388 \pm 56
SSTc2d J163154.7-250324 (SB2) ^e	20/06/2010	GJ 699	-12.1 \pm 1.0	0.56 \pm 0.01	-8.4 \pm 0.2	394 \pm 38
	22/06/2010	GJ 699	-16.4 \pm 1.2	0.5 \pm 0.2	-6 \pm 1	301 \pm 27
	14/06/2011	GJ 514	-5.4 \pm 0.5	0.5 \pm 0.2	-7 \pm 2	397 \pm 39
	mean		-11 \pm 2	0.5 \pm 0.2	-7 \pm 2	364 \pm 61
SSTc2d J163205.5-250236	21/06/2010	GJ 908	-6.7 \pm 0.3	0.50 \pm 0.09	-34 \pm 2	461 \pm 43
	17/01/2012	GJ 908	-6.0 \pm 0.4	0.51 \pm 0.01	-25 \pm 2	427 \pm 39
	mean		-6.4 \pm 0.5	0.51 \pm 0.09	-29 \pm 3	444 \pm 58
SSTc2d J163355.6-244205	22/06/2010	GJ 699	-6.1 \pm 0.4	0.52 \pm 0.02	-2.7 \pm 0.6	255 \pm 23
	14/06/2011	GJ 514	-5.8 \pm 0.3	0.53 \pm 0.02	-7 \pm 3	283 \pm 31
	mean		-6.0 \pm 0.5	0.53 \pm 0.03	-5 \pm 3	269 \pm 39

^aIf the spectral line could not be distinguished from the continuum, a 2σ limit is shown.

^bH α 10% velocity width (White & Basri 2003) variability is due to stellar activity or variable accretion (e.g. Jayawardhana et al. 2003; Natta et al. 2004).

^cThese targets were found to be active stars or variable accretors and are discussed in greater detail in Appendix 8.1.

^dWhite et al. (2007) measured an RV of -1.46 ± 2.21 kms⁻¹, consistent with our measurements.

^eThe RV listed here is the systemic velocity of the binary system, assuming mass ratio $q = 0.95$ (See Section 4). The EWs listed are the blended values at each epoch.

^fNote that SR21A shows H α in absorption rather than in emission. See Appendix 8.2 for further discussion of this object.

Table 3. Radial Velocity Standards

Name	Spectral Type	RV ^a kms ⁻¹
GJ 514	M1	14.56±0.40
GJ 699	M4	-110.51±0.40
GJ 908	M2	-71.15±0.40

^aRVs are taken from Nidever et al. (2002), who quote an RV stability of 0.05 kms⁻¹. The zero point of the absolute RVs in this study is uncertain at the 0.4 kms⁻¹ level (Marcy & Benitz 1989).

3.2. Rotational Velocities

We determined rotational velocities ($v \sin i$) following the ‘Fourier Method’ (e.g. Carroll 1933; Gray 1976; Reiners et al. 2001; Gray 2005; Simón-Díaz & Herrero 2007). Using telluric lines, we determined a small amount of instrumental broadening ($\sim 0.2\text{\AA}$), which we absorbed into model lines of zero rotation. We extracted the excess broadening due to rotation using the relatively high S/N Li I doublet at 6708 \AA and Ca I at 6718 \AA . While forbidden atomic transitions such as [Fe III] do not run the risk of confusing rotational broadening for thermal broadening, their S/N were consistently too low for use.

Our results were sensitive to the assumed limb-darkening coefficient, ϵ , with an $\sim 15 \text{ kms}^{-1}$ deviation in $v \sin i$ for a change $\Delta\epsilon \sim 0.1$. Due to this sensitivity, we used the Claret (2000) linear limb-darkening coefficient catalog for each source, using characteristic values of $\log g$ and effective temperatures (Gizis 1997; Casagrande et al. 2008; Rajpurohit et al. 2013) for their literature SpTs (for G and K stars) and those determined by their TiO-7140 indices (for M stars; see Section 3.3).

For each observation, we took a S/N-weighted average of the $v \sin i$ as measured from the Li I and Ca I lines. For each target, we took the average of the per-observation measurements. Uncertainties were added in quadrature. Two targets in our sample have $v \sin i$ measurements reported in the literature: RX J1852.3-3700 ($23.05 \pm 3.59 \text{ kms}^{-1}$; White et al. 2007) and SSTc2d J162740.3-242204 ($14.7 \pm 0.9 \text{ kms}^{-1}$; Torres et al. 2006). Both values are within the uncertainties of our measured values. We present our measurements in Table 4.

3.3. SpT measurements

M dwarf spectra are dominated by TiO molecular bands. Wilking et al. (2005) defined the TiO-7140 index as a way of determining the SpTs of M dwarfs: a ratio of the mean flux of two 50 \AA bands centered on 7035 \AA (the continuum band) and 7140 \AA (the TiO band). We used the calibration given in Shkolnik et al. (2009)

$$\text{SpT} = (\text{TiO-7140} - 1.0911)/0.1755 \quad (1)$$

to measure the SpTs of all of our targets. Strictly, this relation holds for SpTs M0–M5, which covers the range of SpT literature values of M dwarfs in our sample⁴. Our results

⁴An interesting case is that of target SSTc2d J162309.2-241705, for which C10 report an uncertain spectral type of “M?”. We measure a negative M subclass for this target, suggesting that it may be a late K.

are shown in Figure 3. The error bars for literature values are ± 1 subclass, largely from C10. The error bars for our own measurements are the root mean square=0.6 scatter of the Shkolnik et al. (2009) calibration, added in quadrature with the standard deviation of the measured SpTs across the multiple epochs each target was observed. They favor a linear fit of

$$\text{SpT}_{\text{This work}} = (1.5 \pm 0.4)\text{SpT}_{\text{Lit.}} + (-2 \pm 1) \quad (2)$$

i.e. we measure SpTs consistent with the literature values to $\sim \pm 1$ SpT subclass. Our measured SpT per target is shown in Table 4, along with the literature values (see references in Table 1) and $v \sin i$ (see above). We do not revise K-type SpTs with the exception of SSTc2d J163154.7-250324, which we find to be an SB2 (see Section 4). For this system we report the derived SpTs and $v \sin i$ for each component. The lack of TiO in the spectra of the two G stars in our sample (SSTc2d J162646.4-241160 and SR21A) indicates that they are hotter than M stars, but the difference between very young G and K stars not well-constrained (e.g. Beuther et al. 2014), and therefore we do not revise their SpTs.

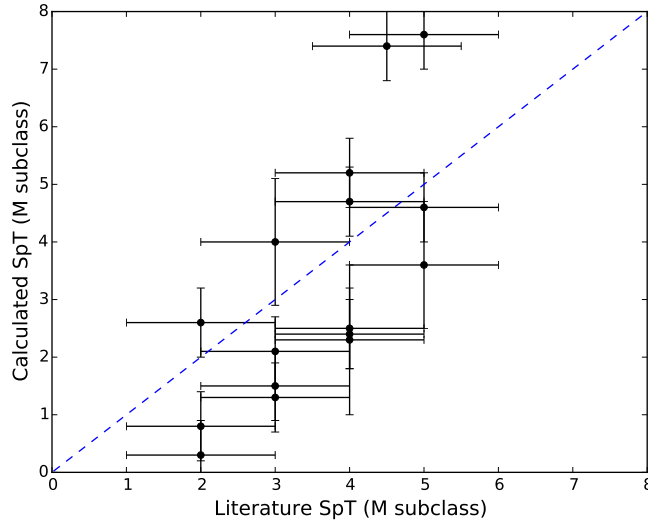


Fig. 3.— The M subclasses of targets reported as M dwarfs in the literature (i.e. $0 \rightarrow M0$, $1 \rightarrow M1$, etc.), with their literature subclasses compared with our own measurements based on their TiO-7140 indices. A 1:1 relation is overlaid in blue. Our measurements are broadly consistent with the literature values. Larger deviations from the literature values at subclasses ≥ 5 occur because the TiO-7140 index is less well-constrained for late M SpTs.

Table 4. Spectral types and rotational velocities

Name	SpT (Lit.)	SpT (This work)	$v \sin i$ kms^{-1}
SZ Cha	K0		34.0 ± 0.4
T25	M2.5	$M1.4 \pm 0.6$	24.3 ± 0.4
T35	K8		29 ± 3
RX J1852.3-3700	K7		19.5 ± 0.5^a
CrA-4111	M4.5	$M7.4 \pm 0.6$	65 ± 6
G-49	M4	$M5.2 \pm 0.6$	43 ± 4
G-102	M5	$M7.6 \pm 0.6$	41 ± 4
SSTc2d J162118.5-225458	M2	$M2.6 \pm 0.6$	40 ± 3
SSTc2d J162218.5-232148	K5		35 ± 16
SSTc2d J162245.4-243124	M3	$M2.1 \pm 0.6$	42.3 ± 0.2
SSTc2d J162309.2-241705	M?	K?	45.0 ± 0.9
SSTc2d J162332.8-225847	M5	$M4.6 \pm 0.6$	49 ± 4
SSTc2d J162336.1-240221	M5	$M3.6 \pm 1.1$	64 ± 5
SSTc2d J162506.9-235050	M3	$M1.3 \pm 0.6$	46 ± 6
SSTc2d J162623.7-244314	K5		10 ± 1
SSTc2d J162646.4-241160	G5		49 ± 5
DoAr28	K?		27 ± 19
SR21A	G3		65 ± 3
SSTc2d J162738.3-235732	K5		42 ± 2
SSTc2d J162739.0-235818	K6		43 ± 23
SSTc2d J162740.3-242204	K5		18 ± 2^b
SSTc2d J162802.6-235504	M3	$M4.0 \pm 1.1$	60 ± 12
SSTc2d J162821.5-242155	M3	$M1.5 \pm 0.6$	63 ± 6
SSTc2d J162854.1-244744	M2	$M0.3 \pm 0.6$	49 ± 7
SSTc2d J163020.0-233108	M4	$M2.4 \pm 0.6$	50 ± 6
SSTc2d J163033.9-242806	M4	$M4.7 \pm 0.6$	60 ± 8
SSTc2d J163145.4-244307	M4	$M2.3 \pm 1.3$	57 ± 12
SSTc2d J163154.4-250349	M4	$M2.5 \pm 0.7$	57 ± 14
SSTc2d J163154.7-250324A ^c	K7	$K7.0 \pm 0.5$	18.2 ± 0.9
SSTc2d J163154.7-250324B ^c		$K9.0 \pm 0.5$	23 ± 5
SSTc2d J163205.5-250236	M2	$M0.8 \pm 0.6$	48 ± 5
SSTc2d J163355.6-244205	K7		43 ± 4

^aWhite et al. (2007) report a $v \sin i = 23.05 \pm 3.59 \text{ kms}^{-1}$ for this target, within 1σ of our measured value.

^bTorres et al. (2006) report a $v \sin i = 14.7 \pm 0.9 \text{ kms}^{-1}$ for this target, within 2σ of our measured value.

^cThis target is an SB2. In this Table we report the properties of each component. See Section 4.

4. Spectroscopic Binary Detection

Single-lined SBs, SB1s are revealed by significant RV variability between observations. We performed χ^2 tests on the RVs of targets with single-peaked cross-correlation functions (CCFs) using the target’s average RV as a flat prior and found no SB1s in our sample of 31 TD objects. As mentioned above, the fact that all of the RVs agree with both the sample average for their SFR and the association average from the literature argues against any of our targets being long-period SB1s, or $P \sim 1$ yr systems whose RV was serendipitously measured by us at the same orbital phase each year.

Cross-correlation of the target spectra with an RV standard spectrum reveals whether or not a star is an SB2 if the orbital phase at the time of observation allows resolvable RV motion such that CCF is double-peaked. We found one SB2 in the sample (SSTc2d J163154.7-250324), which we discuss in the subsections below.

4.1. Measured Properties of SSTc2d J163154.7-250324

The stellar features were blended enough on UT100620 and UT100622 that we could not confirm it as an SB2 using those observations alone. SSTc2d J163154.7-250324 clearly exhibited a double-peaked CCF one year later on UT110614. We measured the RVs of the primary and secondary components to be 16 ± 2 km s^{-1} and -27 ± 2 km s^{-1} , respectively. This gives a systemic RV of $\gamma = -5.4 \pm 0.5$ km s^{-1} based on fluxes in the two CCF peaks (flux ratio of 0.68 ± 0.07 ; see below). We measured values of $v \sin i$ for the primary and secondary as 18.2 ± 0.9 km s^{-1} and 23 ± 5 km s^{-1} , respectively, following the method described in Section 3.2 on the unblended absorption lines seen on UT110614. A section of the stellar spectrum and the CCFs of this SB2 are shown in Figures 4 and 5 respectively.

Since we resolved the two CCF peaks, we were able to estimate spectral types and component masses of the individual stars. Assuming a flux-weighted relation (where f_i is the integrated flux of Gaussian fits to the cross-correlation peaks at R band wavelengths; Figure 5) between component and integrated SpTs (Cruz & Reid 2002; Reid & Cruz 2002; Daemgen et al. 2007) for primary star A and secondary B :

$$\text{SpT}_{\text{int}} = (f_A \text{SpT}_A + f_B \text{SpT}_B) / (f_A + f_B) \quad (3)$$

Using the Kepler’s Third law we related component masses M_n to velocity amplitudes K_n :

$$M_A/M_B = K_A/K_B \quad (4)$$

We imposed a limit on SpT and magnitude (Daemgen et al. 2007; Shkolnik et al. 2010):

$$\Delta R = M_R(\text{SpT}_B) - M_R(\text{SpT}_A) \quad (5)$$

We measure a flux ratio of 0.68 ± 0.07 for the CCF peaks. Thus, we determine that SSTc2d J163154.7-250324 is composed of a K7(± 0.5) and a K9(± 0.5) star, with mass ratio $q \simeq 0.95$. Using Kepler’s Third Law, we calculated the orbital limits to be $a < 0.6$ AU and period $P < 150$ days. Our measurements of this system are summarized in Table 5.

4.2. Comparison to the literature

Before the SB2 SSTc2d J163154.7-250324 was identified as a TD, it had already been measured to have an infrared excess in its SED by 2MASS and exhibit the visible spectral properties of a YSO (Ratzka et al. 2005, who refer to the system as H α 74 and/or ISO-Oph 207). The disk is unresolved in the CHARM2 catalog (Richichi et al. 2005). Padgett et al. (2008) reported the YSO to reside in the most populated region of the Oph SFR. Evans et al. (2009) reported J163154.7-250324 to have an extinction-corrected temperature and a luminosity of 3100 K and $2.5L_\odot$, respectively. It is listed as a YSO candidate star with extinction from dust, but no evidence of local extinction from a surrounding envelope. Cieza et al. (2009), like Ratzka et al. (2005), found no evidence for multiplicity.

C10 identified the system as a TD and determined it to be a K7-type star, in agreement with our measurements above. Their Li I measurements were too low in S/N and therefore unreported, so we cannot compare them to our own measurements of $0.5 \pm 0.2\text{\AA}$. They detected strong emission from the Ca II triplet, providing further evidence that it is indeed a PMS object. Their measurement of the H α 10% velocity-width (470 km s^{-1}) is similar to our own ($364 \pm 61 \text{ km s}^{-1}$), but variation is expected due to variable accretion and the SB2 nature

Table 5. Properties of the SSTc2d J163154.7-250324 system

Property	System	A	B
SpT	K7 ^a	K7.0 \pm 0.5	K9.0 \pm 0.5
RV (km s^{-1})	-5.4 \pm 0.5	16 \pm 2	-27 \pm 2
$v \sin i$ (km s^{-1})		18.2 \pm 0.9	23 \pm 5
Separation (AU)	<0.6		
Period (days)	<150		
Mass ratio	0.95		

^aCieza et al. (2010)

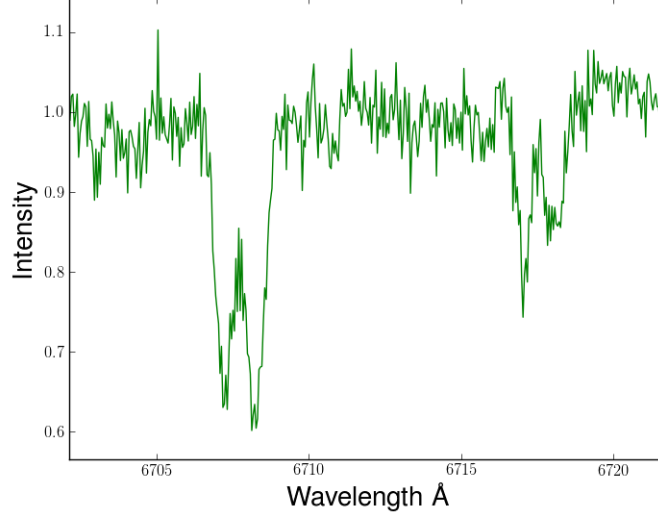


Fig. 4.— The spectrum of SSTc2d J163154.7-250324 on UT110614. The SB2 nature of the system is evident in all absorption features, including the Li lines at 6708 Å.

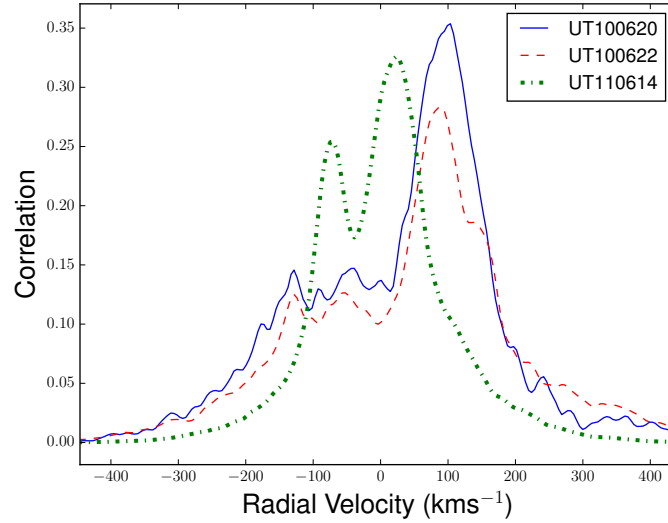


Fig. 5.— The CCFs of SSTc2d J163154.7-250324 averaged across apertures used. Spectra taken in 2010 were cross-correlated against GJ 699 and those in 2011 against GJ 514.

of the object. C10 found no evidence for a companion, as their study was not sensitive to separations < 8 AU. They derive a limit on disk mass of $M_D < 1.1 M_{\text{Jup}}$, and classify it as a grain growth-dominated TD (as mentioned in Section 1, this is their most common TD subclass).

5. Sample sensitivity to companions

We proceeded in a completeness study of our results following Section 6.1 of Duquennoy & Mayor (1991) and Section 3.1 of Melo (2003), considering the observation cadence for each target, $v \sin i$, and the RV uncertainties of each spectrum. We sought to constrain how sensitive we were to SB1s within each TD for different orbital periods and secondary masses.

We estimated the primary mass of each target based on its absolute magnitude using low-mass solar metallicity stellar evolution models from Baraffe et al. (2015). To estimate absolute magnitudes we apply the distance to Oph of 140 pc, for Cha targets we used a distance of 160 pc (Feigelson & Lawson 2004), and for CrA and Cor targets we use a distance of 140 pc (Sicilia-Aguilar et al. 2008). We use the Baraffe et al. (2015) 1 Myr CFHT tracks for all targets (e.g. Furlan et al. 2009; Kim et al. 2009). The 2MASS J magnitudes, available for all of our targets (and K where available), coupled with these age and distance estimates allowed us to estimate the mass of each target. We did not correct for extinction; NIR magnitudes are the least affected by optical- and NIR-excesses from accretion and inner disks. The average change in mass estimate between the Baraffe et al. (1998) and Baraffe et al. (2015) models was $|\Delta M| = 0.2 M_{\odot}$. Such a difference has little impact on our results.

With these estimated primary masses of each target, we generated 10^6 artificial companions, each with mass ratio (q), period (P), time of periastron, inclination (i), and longitude of periastron (ω) chosen randomly from a uniform distribution. Eccentricity (e) was chosen to be zero for $P < 8$ days, and drawn from the Hyades distribution of Burki & Mayor (1986) for $P > 8$ days. The solution of Kepler’s Equation followed Chapter 2.5 of Hilditch (2001). At each date of observation, we calculated the RV for each artificial companion, and added to this a random uncertainty drawn from a Gaussian distribution centered around our measured uncertainty of the (real) primary RV (Table 2).

To avoid an overestimate of detection probability at high mass ratios and long orbital periods (where we expect SB2 systems to lie in this $(M_{\text{secondary}}, P)$ -space) we calculated the flux-weighted systemic velocity of the binary system

$$\gamma = \frac{\text{RV}_1 A_1 + \text{RV}_2 A_2}{A_1 + A_2} \quad (6)$$

where A_2/A_1 is the flux ratio of the stars, and RV_1 and RV_2 are the RVs of each star. Their orbits are related by angle $\theta_2 = \theta_1 + \pi$, orbiting the center of mass. We interpolated the fluxes in the I-band using the Baraffe et al. (2015) 1 Myr CFHT tracks. Defining $\alpha_1 = \cos(\theta_1 + \omega) + e \cos \omega$, $\alpha_2 = \cos(\theta_1 + \pi + \omega) + e \cos \omega$, and expressing the flux ratio as $F_R = A_2/A_1$, the systemic RV of the system is then

$$\gamma = \frac{2\pi a \sin i}{(1 + F_R)P(1 - e^2)^{1/2}}(M_2\alpha_1 + M_1\alpha_2F_R) \quad (7)$$

We then tested if the artificial companion could be ‘detected’ by whether or not the root-mean-square of the values of γ (each given by Equation 7) was greater than 0.7 kms^{-1} , i.e. greater than the upper limit on the standard errors of our non-SB RVs. With over 10^6 simulations per target, we were able to map probabilities of detection in $(M_{\text{secondary}}, P)$ -space. The results of these simulations are shown in Figure 6, which shows contour levels in 10% bins from 40 – 100% chance of detection⁵. The periodic modulation in the contour patterns demonstrates our insensitivity to orbital periods that are multiples of our observing cadence.

The results of these simulations allow us to accurately assess our sensitivity to companions on a per-target basis, with a degeneracy in P and mass ratio q .⁶

For example, for SZ Cha we are sensitive to ($\geq 80\%$ chance of detection) a $0.4 M_{\odot}$ companion at periods $P < 30$ days and $70 < P < 75$ days, whereas for SSTc2d J162623.7-244314, we are sensitive to a $0.4 M_{\odot}$ companion at all periods. Such an example illustrates the challenge in quoting an “overall” or “average” sensitivity for this non-homogeneous sample. We are, however, consistently sensitive to short-period ($P < \sim 100$ days) SBs for all of our targets.

6. Discussion

Prato (2007) measures a spectroscopic binary fraction of $0.12^{+0.08}_{-0.04}$ among Oph K7-M4 stars. Of the 33 targets in Prato (2007), two are TD objects but neither appears to be an SB: ROXR1 20 (C10) and RX J1614.4-1857 (Wahhaj et al. 2010). Of the five SBs they found, one hosts a debris disk (RX J1612.6-1924, Wahhaj et al. 2010) and the rest are diskless (RX

⁵While contour levels run between these values, they rarely go below a 60% chance of detection. This happens only for the lowest secondary masses.

⁶One can convert between limits on the period and limits on the semi-major axis of separation a using the relation: $a(P, q, M_{\text{primary}}) = (P/2\pi)^{2/3}(G(1 + q)M_{\text{primary}})^{1/3}$

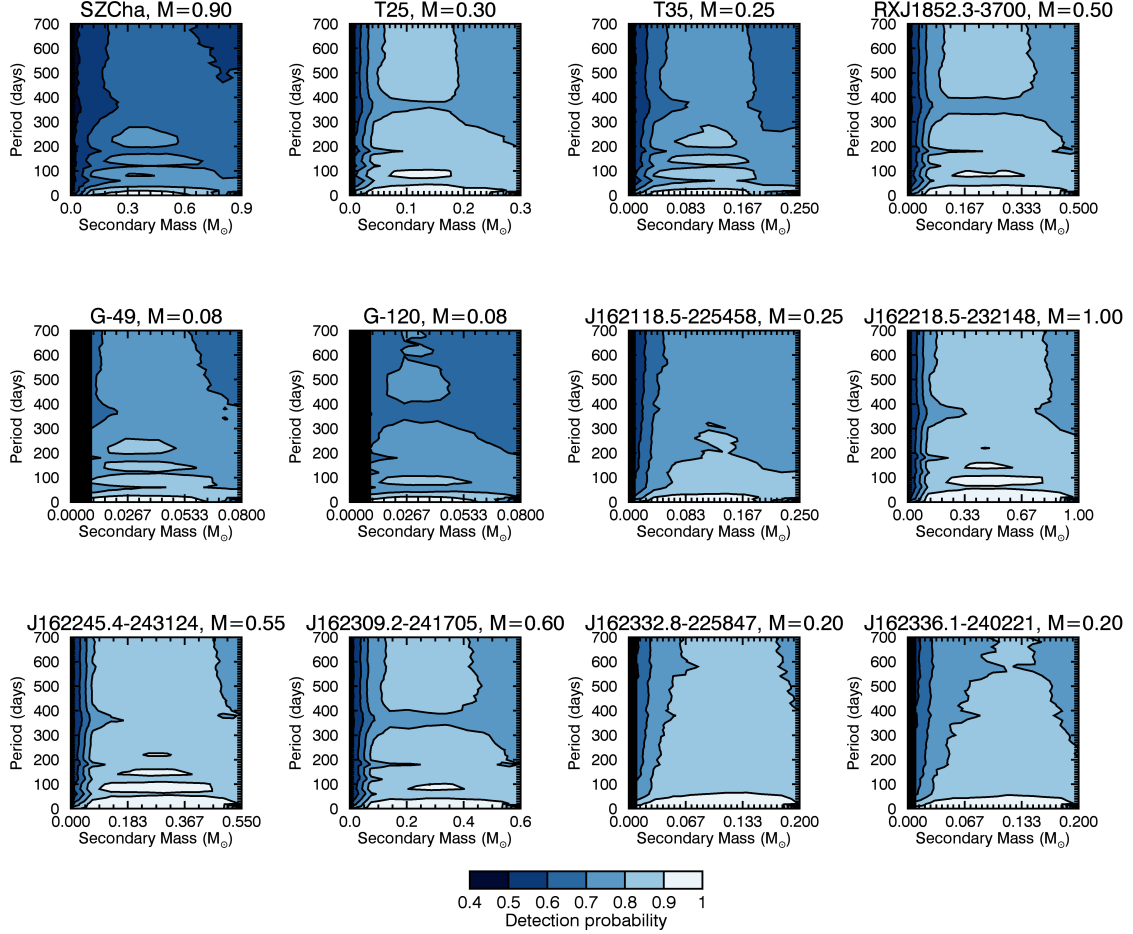


Fig. 6.— Calculated contours in $(M_{\text{secondary}}, P)$ -space representing the probability of detecting a companion star spectroscopically, given the precision and timing of our RV measurements. Contour levels of probability of detection are shown in 10% bins from 40 – 100%. Shown next to the name of the target system is the mass of the primary in solar masses, according to the PMS models of Baraffe et al. (2015). We imposed that the secondary mass not exceed the primary mass for each system. Two targets are not shown since we have only one RV measurement for the star, and therefore no chance of finding SB1s: these are SR21A and CrA-4111. For the Ophiuchus targets, we have dropped the ‘SSTc2d’ prefix.

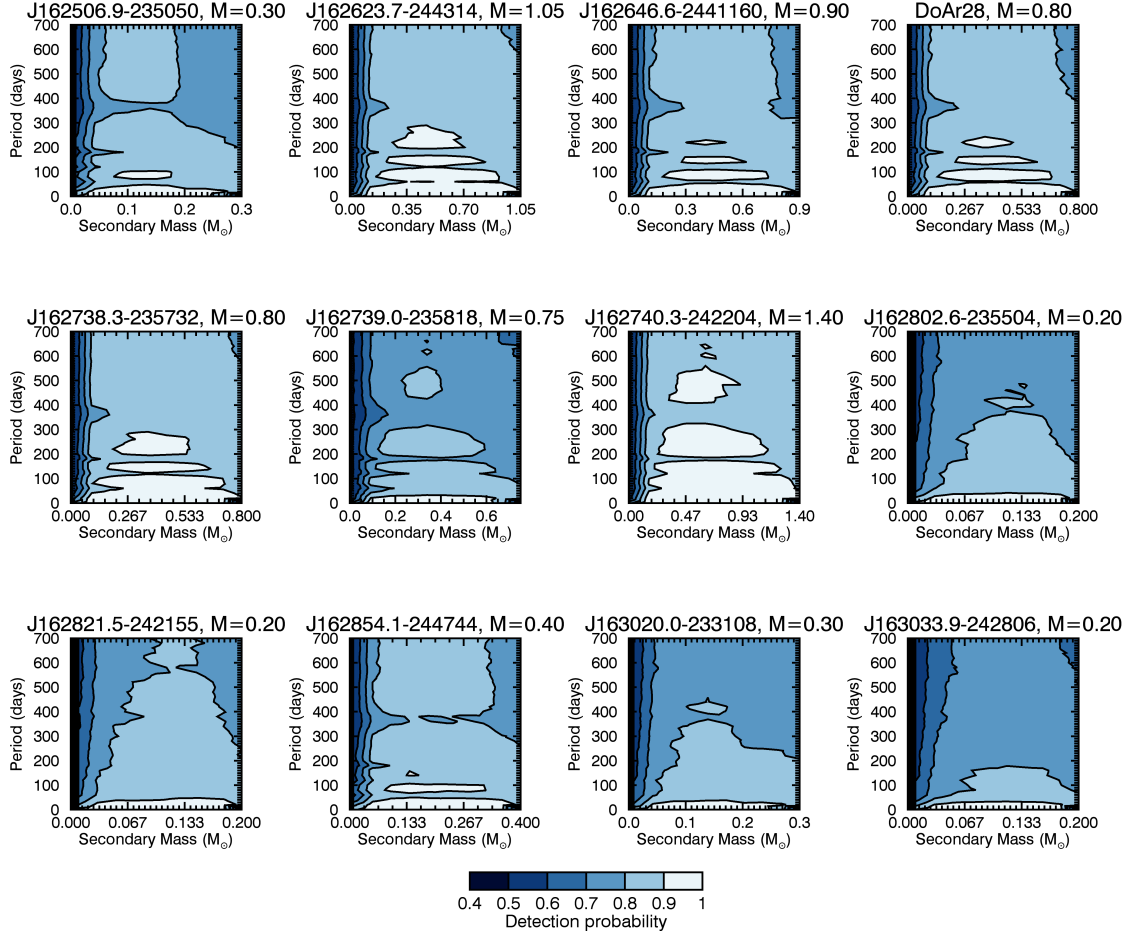


Fig. 6.— (continued.)

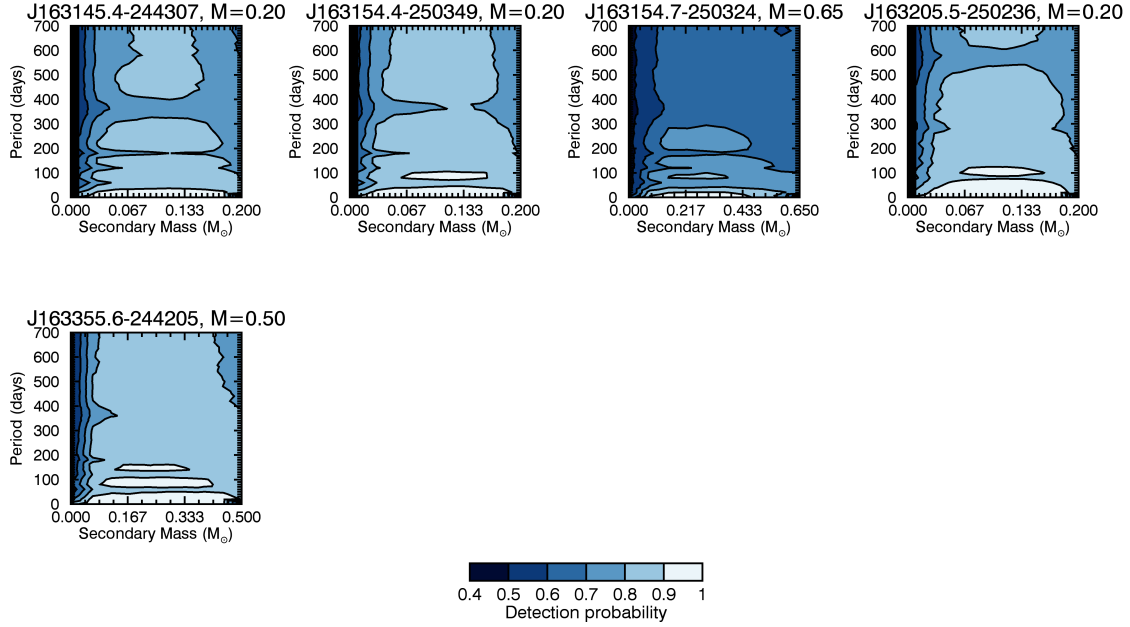


Fig. 6.— (continued.)

J1612.3-1909, RX J1622.7-2325, and RX J1622.8-2333; Wahhaj et al. 2010; ROXR1 14 Cieza et al. 2007; Rosero et al. 2011). We found one SB2 in a sample of 24 Oph TD objects, and no SBs among the TD objects in the other SFRs surveyed. Including the two TD targets from Prato (2007) in our Oph sample, we find that multiplicity among TD objects in this region to be 1/26. We measure an SB fraction of $0.04^{+0.12}_{-0.03}$, determining uncertainties following the binomial theorem. This is consistent with that of the non-TD, late-type stars in Ophiuchus (Prato 2007) and the young stars in Chameleon and Taurus-Auriga ($0.07^{+0.05}_{-0.03}$ and $0.06^{+0.03}_{-0.02}$, respectively; mass range $0.2-3 M_{\odot}$; Nguyen et al. 2009). The result also in agreement with Lodieu et al. (2014), who report an SB fraction of 0.054 ± 0.038 for wide binaries (50 – 150 AU) in which one star is a planet-host. Their SB fraction for wide, planet-hosting substellar binaries is also within our range at 0.027 ± 0.027 . We also conclude that there is no significant difference between our result and that among the low-mass ($0.08 M_{\odot}-0.6 M_{\odot}$) field stars in the surveys of Duquennoy & Mayor (1991) and Fischer & Marcy (1992), $0.09^{+0.03}_{-0.02}$ and $0.03^{+0.04}_{-0.02}$, respectively. Likewise, Raghavan et al. (2010) found an SB fraction of $0.073^{+0.014}_{-0.012}$ among nearby old solar-type dwarfs.

That the fraction of SBs in TDs is similar to diskless stars (within uncertainties) may suggest that whatever causes the disk dissipation is independent of binarity, i.e. disks around close SBs evolve on a similar timescale as those around single stars. Such a conclusion would mean that planet formation timescales around SBs may be comparable to those around single

stars.

However, several studies (Ghez et al. 1997; White & Ghez 2001; Cieza et al. 2009; Duchêne 2010; Kraus et al. 2012) have found protoplanetary disk fractions are lower for young (≤ 5 Myr) binary systems with $a < 50$ AU than for single stars or wider binaries. With the exception of Kraus et al. (2012), these studies were not sensitive to the tight binary systems to which our study was especially sensitive. Kraus et al. (2012) and Cheetham et al. (2015), sensitive to similar binary separations to our study, find the disk fraction among close visual binaries (0.1–40 AU) in Taurus-Auriga and Ophiuchus to be lower than for single stars. Their studies did not distinguish between circumbinary TDs and other disk types (e.g. debris disks). If the SB fraction of TD systems is actually lower than that in the diskless population, this suggests that the binary TD stage of a close-in binary is even shorter than the TD stage of single stars.

While we have searched for SBs specifically in TDs, Prato (2007) made a complementary search for SBs irrespective of knowledge of their disks. None of the SBs she found were in transitional or protoplanetary disks. This is consistent with SBs having a faster TD phase and that the overall disk fraction among SBs remains low down to small separations.

Such a conclusion would be in contention with the results of Alexander (2012), whose models suggest that protoplanetary disks around tight binaries ($a \leq 1$ AU) are longer-lived than those around wider binaries. This is because the very close binaries efficiently clear the inner disk to radii larger the characteristic radius of photoevaporative winds, suppressing total disk dispersal.

That we only find one SB in our TD sample could reflect that dissipation is so fast among these objects that most disks are completely dispersed by the age of Oph. However, we do not have the statistics to draw a steadfast conclusion. But since we find no evidence that TD objects have higher SB fractions than other stellar populations, we can infer that close-in binaries are not likely the primary cause of inner holes in TDs. This leads to giant planet formation as a more likely explanation.

7. Summary

We have presented a spectroscopic survey of 31 TD stars, 24 of which lie in Ophiuchus. We found one of the Oph stars to be an SB2 (SSTc2d J163154.7-250324). This system is composed of a K7(± 0.5) and a K9(± 0.5) star, with mass ratio $q \simeq 0.95$ and orbital limits $a < 0.6$ AU and $P < 150$ days. The average RV of our Oph targets is -6.6 kms^{-1} , with a dispersion of 1.3 kms^{-1} . The median uncertainty in the measured RVs including systematic

uncertainties is 0.4 km s^{-1} . The average RVs of all four clouds surveyed (Cha, Cor, CrA and Oph) are consistent with the literature values. With this single SB detection, we measure an SB fraction of $0.04^{+0.12}_{-0.03}$. This finding is consistent with that of non-TD late-type stars in and outside of the region. This suggests that a TD may not be dispersed more efficiently by a tight binary than by a single star, and may imply that planet formation timescales around close binaries are comparable to those around single stars.

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Facility: LCO: Magellan (MIKE).

8. Appendix

We present here the accretion properties and $\text{H}\alpha$ and Li I measurements of individual stars.

8.1. Age and accretion diagnostics

The strength of the Li I line ($\lambda=6708 \text{ \AA}$) is a standard diagnostic for the age of low-mass stars, since lithium is rapidly depleted in their atmospheres. All but one of our targets exhibit Li I absorption, so we are only able to place upper limits on the absorption in the spectra of CrA-4111 due to low S/N in our single observation of the system. The Li I should be undetectable in early M stars after $\sim 20 \text{ Myr}$ (Baraffe et al. 1998; White & Hillenbrand 2005). For G and K stars the depletion can take much longer ($\geq 300 \text{ Myr}$, e.g. Zickgraf et al. 2005). The age of Oph is $< 10 \text{ Myr}$ (Duncan 1981; Chabrier et al. 1996; Martin et al. 1999a,b; Zuckerman & Song 2004, C10), and thus, all stars are expected to display Li absorption. For the 12 stars at sufficient S/N to detect Li that were also observed in C10, our EWs are consistent with theirs. Our Li I EW measurements of DoAr 28 and SR21A are consistent

with those of Magazzu et al. (1992) and James et al. (2006), respectively. We are not aware of any previous measurements of Li I in SSTc2d J162309.2-241705. The Li I EWs are listed in Table 2.

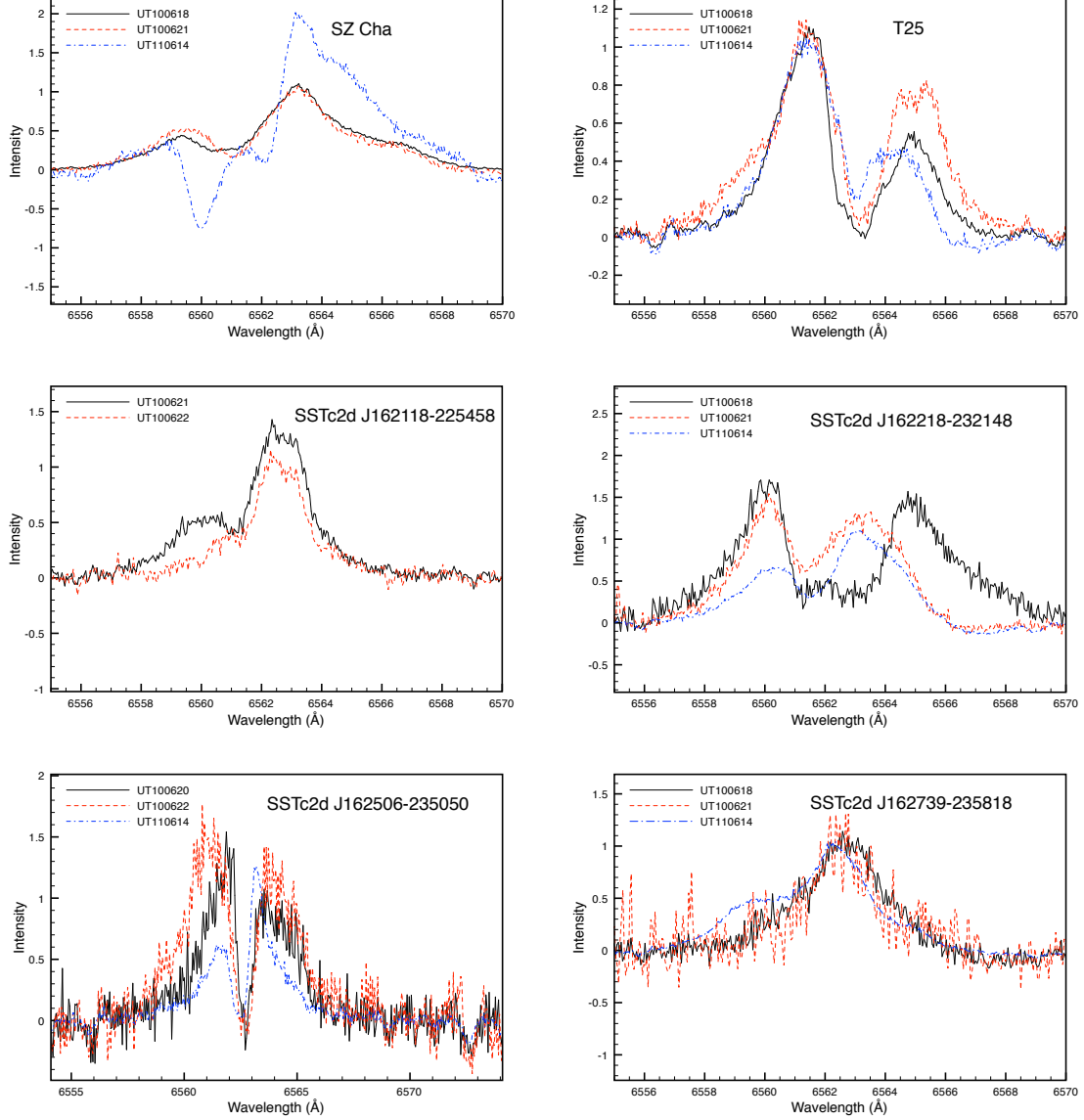
H α ($\lambda=6563$ Å) is an indicator of stellar activity and gas accretion onto a star. Applying the relation found by Natta et al. (2004), we use the 10% velocity-width (White & Basri 2003) of the H α emission to estimate the rate of accretion from the disk onto the star for targets with H α 10% velocity-widths greater than 200 kmbs⁻¹ (this is true for 21 of our 31 targets). We find that of the 15 targets that were observed by C10 and had sufficient S/N at H α , 10 have 10% velocity-widths within 3σ of those measured by C10. The large variations in the remaining measurements are likely due to variable accretion upon the star (e.g. Nguyen et al. 2009). Six of our targets exhibited varying H α during our observations as shown in Figure 7. These are SZ Cha, T25, SSTc2d J162118-225458, SSTc2d J162218-232148, SSTc2d J162739-235818, and SSTc2d J162506-235050. The level of variability is unique for each target, but on average the EW varies by $\sim 80\%$.

8.2. SR21A

At about 1–3 Myr old (which we note is older than the average age of Oph stars; Siess et al. 2000; Andrews et al. 2009; Brown et al. 2009; Follette et al. 2013), SR21A is the primary component of a wide binary system in Oph (Barsony & Marsh 2005). Its spectral type has been reported to be an early G (Suarez et al. 2006; Andrews & Williams 2007b), an early K (Struve & Rudkjobing 1949), and an M star (Kimeswenger et al. 2004). Our measurements of its TiO-7140 index (see Section 3.3) suggest it is not an M-type star ($\text{SpT}_{\text{TiO}} < 0$). SR21A displays H α in absorption rather than in emission,⁷ the only such star in our sample. Li I absorption is also weaker than all of our other Oph objects. The H α absorption could be due to variable accretion (Johns-Krull & Valenti 2001; Baraffe et al. 2012; Hamilton et al. 2012). This is suggested by the broadened blue shoulder of the absorption (a feature possibly hidden by the N II absorption on the red shoulder), a sign of possible accretion emission superimposing upon the photosphere absorption (Figure 8, left). Weak Li I absorption (0.13 ± 0.02 Å; Figure 8, right) might suggest that SR21A is older than expected (>10 Myr; Chabrier et al. 1996), but the existence of a transition disk is evidence against the star being >15 Myr (Mamajek 2009; Muzerolle et al. 2010). While James et al. (2006) find a slightly

⁷It should be noted that SR21A is listed as an emission star in the SIMBAD database, but we did not observe it as such. Struve & Rudkjobing (1949) do list it as a non-emitting star in their spectral emission study of Oph.

Fig. 7.— Variable accretion shown in $H\alpha$.



higher EW for Li I in SR21A (0.275 ± 0.028 Å, versus our value of 0.13 ± 0.02 Å), they also note its inconsistency with other EWs in Oph. It is possible that the same episodic accretion that could be causing the $H\alpha$ absorption has also resulted in abnormal Li depletion in this star (Baraffe et al. 2009; Baraffe & Chabrier 2010; Baraffe et al. 2012). Depletion is also more realistic than the feature being veiled by accretion, since a disk hot enough for veiling

would have caused SR21A to be disqualified as a TD object. However, we compared the EWs/line strengths of all of the lines in the same echelle order as the Li line in SR21A with the late K- and early M-type stars we observed on the same night (UT110614). We found that SR21A had weaker absorption for all lines, suggesting that veiling may still play a role. Overall weaker lines might instead indicate a lower overall metallicity compared with the other Oph targets.

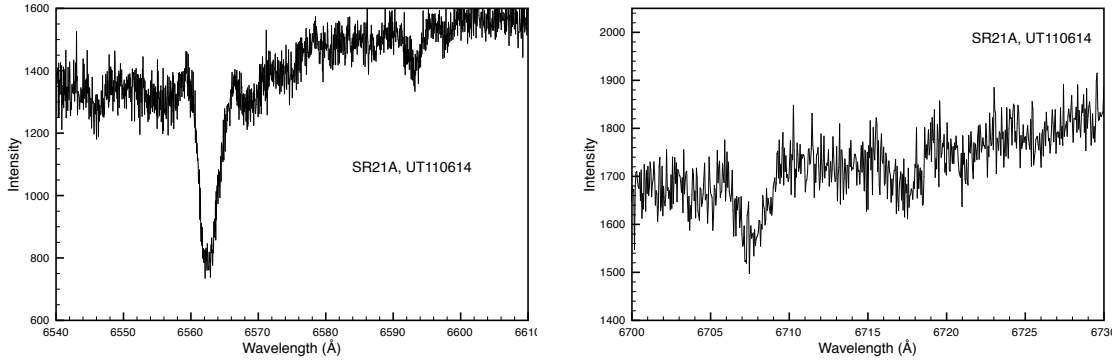


Fig. 8.— Spectral features of SR21A: $H\alpha$ absorption (*left*; $EW=1.13\pm0.15$ Å) and weak Li I absorption (*right*; $EW=0.13\pm0.02$ Å).

An extensive study of the disk morphology of SR21A was recently presented by Follette et al. (2013). Their results may support the postulate of a substellar companion present within the SR21A disk at ~ 18 AU from the primary (Eisner et al. 2009). Since we have just one observation of this system, we can neither confirm nor exclude this finding.

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